

## Nuclear magnetic resonance effect and adaptation

Padmalochan Panda, Uma Parida, Jitendra Kumar Kar, Subrat Kumar Biswal

Department of Physics, NM Institute of Engineering and Technology, Bhubaneswar, Odisha

Department of Physics, Raajdhani Engineering College, Bhubaneswar, Odisha

Department of Physics, Aryan Institute of Engineering and Technology Bhubaneswar, Odisha

Department of Basic Science & Humanities, Capital Engineering College, Bhubaneswar, Odisha

**ABSTRACT:** A prototypical electron-doped iron-based superconductor  $\text{LaFeAsO}_{1-x}\text{H}_x$  goes through an antiferromagnetic (AF) stage for  $x \geq 0.49$ . We have performed atomic attractive reverberation (NMR) estimations on  $\text{LaFeAsO}_{0.4}\text{H}_{0.6}$  at 3.7 GPa to research the attractive prop-erties in the region of a pressing factor incited quantum basic point (QCP). The linewidth of  $1\text{H-NMR}$  spectra expands at low temperatures under 30 K, proposing that the turn minutes stay requested at 3.7 GPa. The conjunction of gapped and gapless turn ex-references was affirmed in the arranged state from the unwinding time  $T_1$  of 75As. The pressing factor initiated QCP is assessed to be 4.1 GPa from the pressing factor reliance of the gapped excitation.

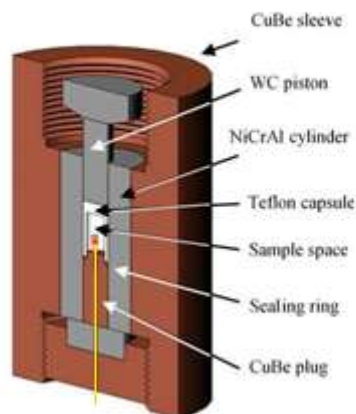
### I. INTRODUCTION

A prototypical electron-doped iron-based pnictide  $\text{LaFeAsO}_{1-x}\text{H}_x$  ( $0 < x < 0.6$ ) exhibits unique electronic properties in a heavily carrier-doped regime: a superconducting (SC) phase with double-domes structure expands in a wide regime ( $0.05 < x < 0.49$ ) [1] and an antiferromagnetic (AF) phase manifests itself by further H doping ( $0.49 < x$ ) [2–4]. Band calculations show that both Fermi surfaces and nesting vectors change by H doping: the two hole pockets present at  $\Gamma$  point in the lightly H-doped regime almost disappear in the heavily H-doped regime [5,6]. The change in the nesting vectors due to H doping would cause each change in wave-vector ( $q$ ) dependent spin susceptibility  $\chi(q, \omega)$  and would allow for the appearance of two AF phases in the lightly and heavily H-doped regimes.

The AF phase in the heavily H-doped regime is strongly suppressed upon applying pressure [7]. We have performed nuclear magnetic resonance (NMR) measurements on  $\text{LaFeAsO}_{0.4}\text{H}_{0.6}$  at 3.7 GPa, and we have found that the spin excitation gap appearing at the AF phase vanishes at around 4.1 GPa. We have investigated the magnetic properties in the vicinity of a pressure-induced quantum critical point (QCP) (c. 4.1 GPa).

### II. EXPERIMENTAL APPARATUSES AND CONDITIONS

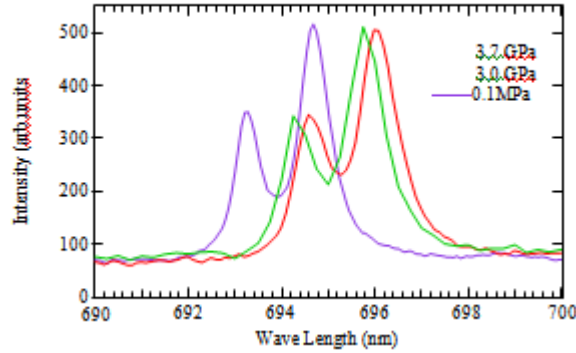
A pressure of 3.7 GPa was applied using a NiCrAl-hybrid clamp-type pressure cell as shown in Fig. 1 [8]. We have used a mixture of Fluorinert FC-70 and FC-77 as the pressure-transmitting medium. A coil wound around the powder samples and an optical fiber with the Ru by powder were inserted into the sample space.



**Figure 1:** A NiCrAl hybrid clamp-type pressure cell [8]. A coil wound around the powder samples and an optical fiber with the Ru by powder were inserted into the sample space.

were inserted into the sample space of the pressure cell [8]. The size of the coil was 2.4 mm in diameter and 3.5 mm in length, and the number of windings was 18 turns. The pressure was monitored through Ruby fluorescence measurements. The R1 and R2 lines at ambient pressure, 3.0 and 3.7 GPa are shown in Fig. 2. The wavelength of the R1 or R2 peak shifts linearly with respect to pressure. The shift of the wavelength  $\Delta\lambda$  satisfies the relation  $P(\text{GPa}) = \Delta\lambda(\text{nm})/0.365$ .

NMR measurements for the powder samples were acquired using a conventional coherent-pulsed NMR spectrometer. The relaxation rate ( $1/T_1$ ) was measured using a conventional saturation-recovery method for the samples whose FeAs planes are parallel to the applied field.



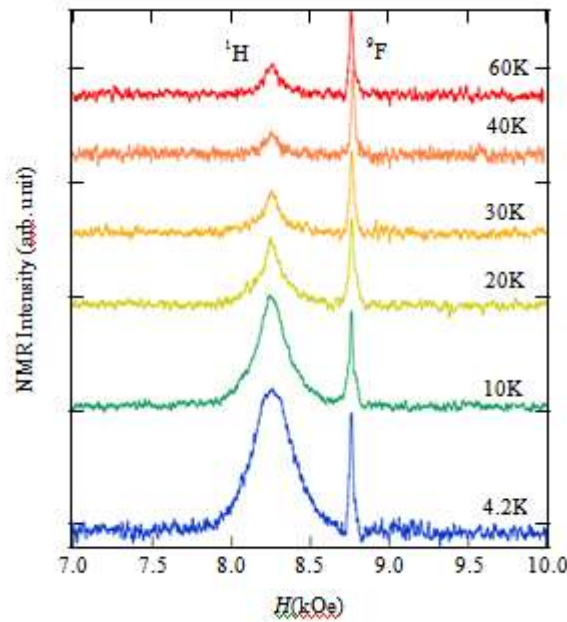
**Figure 2:** Ruby fluorescence spectra. The smaller and larger peaks correspond to the R2 and R1 transitions, respectively.

### III. EXPERIMENTAL RESULTS

#### i. $^1\text{H}$ -NMR spectra

$^{75}\text{As}$  ( $I=3/2$ )-NMR spectra broadened due to the nu-

clear quadrupole interaction, which makes difficult to investigate the antiferromagnetic (AF) state. However,  $^1\text{H}$  ( $I=1/2$ ) is free from the nuclear quadrupole interaction. Therefore, the  $^1\text{H}$  signal is narrow at a paramagnetic state, and the broadening in the AF phase directly reflects the mag-



**Figure 3:**  $^1\text{H}$ -NMR spectra for  $\text{LaFeAsO}_{0.4}\text{H}_{0.6}$  measured at 3.7 GPa and 35.1 MHz. The  $^9\text{F}$  signal originates from the pressure-transmitting medium, a mixture of Fluorinert FC-70 and FC-77.

Figure 4: The increase in  $^1\text{H}$  linewidth due to the ordered spin moments.  $T_N$  represents the antiferromagnetic (AF) transition temperature.

nitude of the spin moments. Figure 3 shows  $^1\text{H}$ -NMR spectra measured at 3.7 GPa and 35.1 MHz. The sharp signal of  $^9\text{F}$  originates from the pressure-transmitting medium mentioned above. The temperature dependence of the linewidth is shown in Fig. 4 together with the data at ambient pressure [2, 4]. The onset of the broadening in Fig. 4 corresponds to the AF transition temperature ( $T_N$ ). The maximum spin moment is estimated to be  $1.80 \mu_B$  [4]. As seen in Fig. 4,  $T_N$  is about 100 K at ambient pressure and decreases to 30 K at 3.7 GPa. The pressure-induced QCP is expected at a much higher pressure regime.

ii.  $1/T_1T$  for  $^{75}\text{As}$

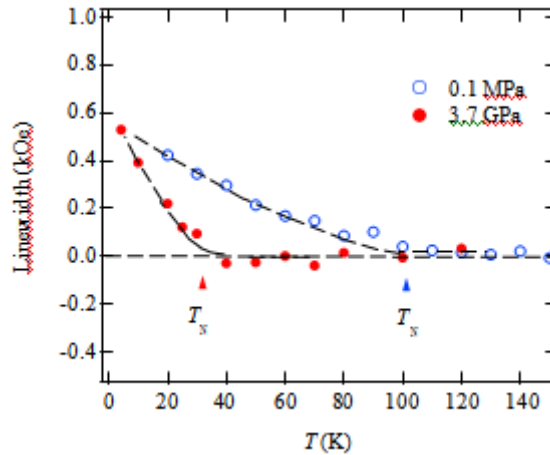


Figure 4: The increase in  $^1\text{H}$  linewidth due to the ordered spin moments.  $T_N$  represents the antiferromagnetic (AF) transition temperature.

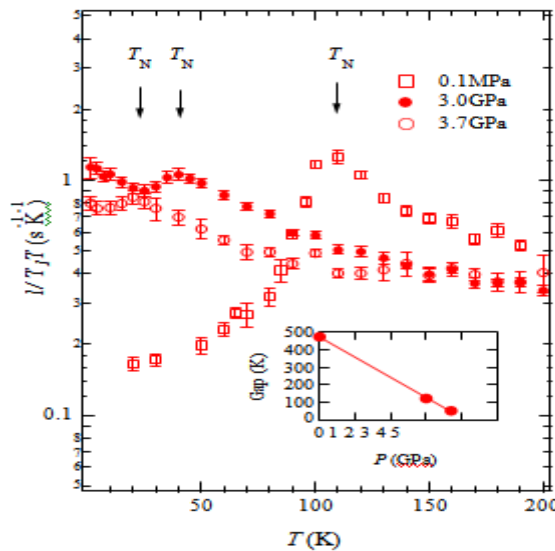


Figure 5: Relaxation rate of  $^{75}\text{As}$  divided by temperature,  $1/T_1T$  for  $\text{LaFeAsO}_{0.4}\text{H}_{0.6}$ .  $T_N$  represents the AF transition temperature. The inset shows the pressure dependence of the spin excitation gap  $\Delta$  (See Eq. (1)).

shows  $1/T_1T$  for  $^{75}\text{As}$ , and the peaks correspond to  $T_N$ . The values of  $T_N$  determined from  $1/T_1T$  are consistent with those obtained from the linewidth of  $^1\text{H}$ . At low temperatures just below  $T_N$ ,  $1/T_1T$  is expressed as follows:

The relaxation rate divided by temperature  $1/T_1T$  provides a measure of low-energy spin fluctuations.

1

$1/T_1T \propto$

$$e^{-T}$$

(1)

In general, neglecting the wave-number ( $q$ ) dependence of the hyperfine coupling constant,  $1/T1T$  is proportional to the imaginary part of the susceptibility:  $1/T1T \propto \sum_q \text{Im}\chi(q, \omega)/\omega$  where  $\omega$  represents a NMR frequency.  $^{75}\text{As}$  is preferred to  $^1\text{H}$  for  $T1$  measurements, because FeAs layers are hardly affected by the random distribution of hydrogen in  $\text{LaO}_{1-x}\text{H}_x$  layers. Furthermore, owing to the nuclear quadrupole interaction, one can pick up the  $^{75}\text{As}$  signals coming from the powders whose FeAs planes are parallel to the applied field. Figure 5 where  $\Delta$  represent the spin excitation gap. The pressure dependence of  $\Delta$  is shown in the inset to Fig. 5. Assuming that  $\Delta$  shows the linear dependence, the pressure-induced QCP is estimated to be 4.1 GPa.

#### IV. DISCUSSION

The activated spin excitation as shown in Eq. (1) originates from a spin density wave (SDW). However,  $1/T1T$  also shows Curie-Weiss behavior below  $T_N$ . The behavior is not observed at ambient pressure and it is characteristic of the critical behavior near the pressure-induced QCP. The coexistence of the gapped and gapless excitations are specific to this system. In this system, major Fermi surfaces are electron pockets with a square-like shape in two-dimensional  $k$ -space. Some parts of the electron pockets would contribute to the nesting and the SDW formation. The critical behavior would originate from the other parts of the Fermi surfaces. The nesting condition becomes worse and the bandwidth becomes broader with increasing pressure. Owing to these effects, the activated behavior shown in Eq. (1) would disappear at the pressure-induced QCP.

#### V. CONCLUSIONS

We performed NMR measurements on  $\text{LaFeAsO}_{0.4}\text{H}_{0.6}$  at 3.7 GPa to investigate the magnetic properties in the vicinity of the pressure-induced QCP. We have found that the SDW ordered state still remains at 3.7 GPa. The pressure-induced QCP is estimated to be 4.1 GPa from the pressure dependence of the spin excitation gap. The gapless excitation observed as the Curie-Weiss behavior of  $1/T1T$  coexists with the gapped excitation, implying that each excitation originates from different parts within the Fermi surfaces.

#### REFERENCE

- [1] S Iimura, S Matsuishi, H Sato, T Hanna, Y Muraba, S W Kim, J E Kim, M Takata, H Hosono, Two-dome structure in electron-doped iron arsenide superconductors, *Nat. Commun.* **63**, 943(2012).
- [2] N Fujiwara, S Tsutsumi, S Iimura, S Matsuishi, H Hosono, Y Yamakawa, H Kontani, Detection of antiferromagnetic ordering in heavily doped  $\text{LaFeAsO}_{1-x}\text{H}_x$  pnictide superconductors using nuclear-magnetic-resonance techniques, *Phys. Rev. Lett.* **111**, 097002(2013).
- [3] M Hiraishi, S Iimura, K M Kojima, J Yamaura, H Hiraka, K Ikeda, P Miao, Y Ishikawa, S Torii, M Miyazaki, I Yamauchi, A Koda, K Ishii, M Yoshida, J Mizuki, R Kadono, R Kumai, T Kamiyama, T Otomo, Y Murakami, S Matsuishi, H Hosono, Introduction to solid state physics, *Nat. Phys.* **10**, 300(2014).
- [4] R Sakurai, N Fujiwara, N Kawaguchi, Y Yamakawa, H Kontani, S Iimura, S Matsuishi, H Hosono, Quantum critical behavior in heavily doped  $\text{LaFeAsO}_{1-x}\text{H}_x$  pnictide superconductors analyzed using nuclear magnetic resonance, *Phys. Rev. B* **91**, 064509(2015).
- [5] Y Yamakawa, S Onari, H Kontani, N Fujiwara, S Iimura, H Hosono, Phase diagram and superconducting states in  $\text{LaFeAsO}_{1-x}\text{H}_x$  based on the multi-orbital extended Hubbard model, *Phys. Rev. B* **88**, 041106(R)(2013).
- [6] S Iimura, S Matsuishi, M Miyakawa, T Taniguchi, K Suzuki, H Usui, K Kuroki, R Kajimoto, M Nakamura, Y Inamura, K Ikeuchi, S Ji, H Hosono, Switching of intra-orbital spin excitations in electron-doped iron pnictide superconductors, *Phys. Rev. B* **88**, 060501(R) (2013).
- [7] N Fujiwara, N Kawaguchi, S Iimura, S Matsuishi, H Hosono, Quantum phase transition under pressure in a heavily hydrogen-doped iron-based superconductor  $\text{LaFeAsO}$ , *Phys. Rev. B* **96**, 140507(R)(2017).
- [8] N Fujiwara, T Matsumoto, K K Nakazawa, A Hisada, Y Uwatoko, Fabrication and efficiency evaluation of a hybrid NiCrAl pressure cell upto 4 GPa, *Rev. Sci. Instrum.* **78**, 073905 (2007).